

بِسْمِ اللَّهِ الرَّحْمَنِ الرَّحِيمِ



سیگنال‌ها و سیستم‌ها

درس ۱۳

تبدیل فوریه‌ی پیوسته-زمان (۱)

The Continuous-Time Fourier Transform (1)

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<http://courses.fouladi.ir/sigsys>

طرح درس

COURSE OUTLINE

استخراج جفت تبدیل فوریه‌ی پیوسته-زمان

Derivation of the CT Fourier Transform pair

مثال‌هایی از تبدیل‌های فوریه

Examples of Fourier Transforms

تبدیل‌های فوریه‌ی سیگنال‌های متناوب

Fourier Transforms of Periodic Signals

خصوصیات تبدیل فوریه‌ی پیوسته-زمان

Properties of the CT Fourier Transform

تبدیل فوریه‌ی پیوسته-زمان (۱)

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استخراج
جفت
تبدیل فوریه‌ی
پیوسته-زمان

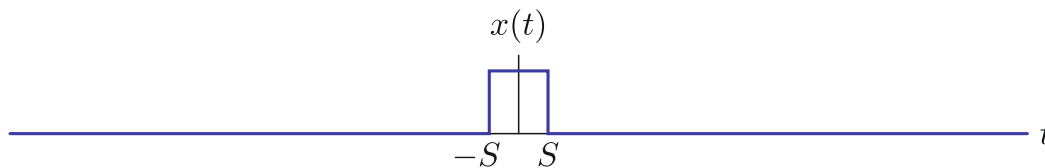
تبدیل فوریه

سیگنال نامتناوب

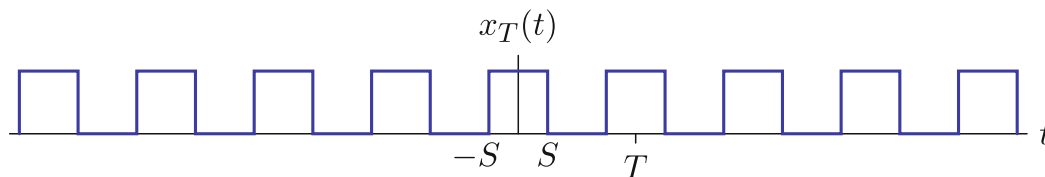
FOURIER TRANSFORM

یک سیگنال نامتناوب می‌تواند به عنوان یک سیگنال متناوب با دوره‌ی تناوب نامتناهی دیده شود.

Let $x(t)$ represent an aperiodic signal.



“Periodic extension”:
$$x_T(t) = \sum_{k=-\infty}^{\infty} x(t + kT)$$

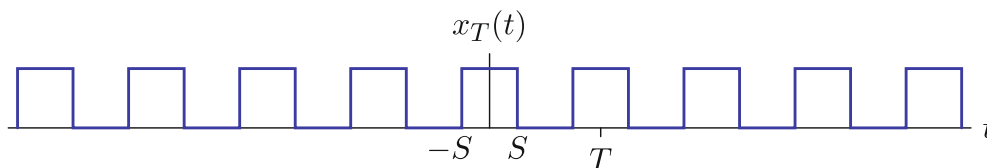


Then
$$x(t) = \lim_{T \rightarrow \infty} x_T(t).$$

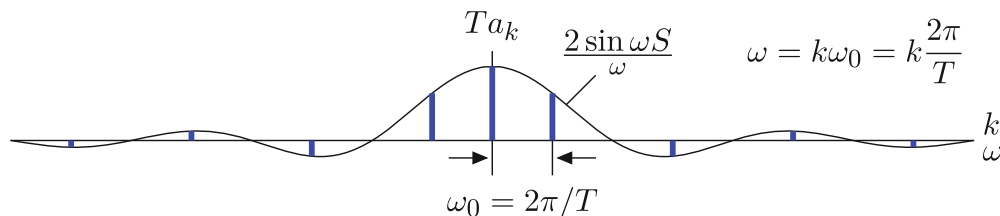
تبدیل فوریه

FOURIER TRANSFORM

Represent $x_T(t)$ by its Fourier series.



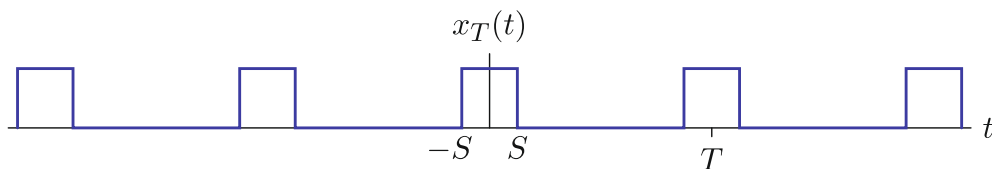
$$a_k = \frac{1}{T} \int_{-T/2}^{T/2} x_T(t) e^{-j\frac{2\pi}{T}kt} dt = \frac{1}{T} \int_{-S}^S e^{-j\frac{2\pi}{T}kt} dt = \frac{\sin \frac{2\pi kS}{T}}{\pi k} = \frac{2}{T} \frac{\sin \omega S}{\omega}$$



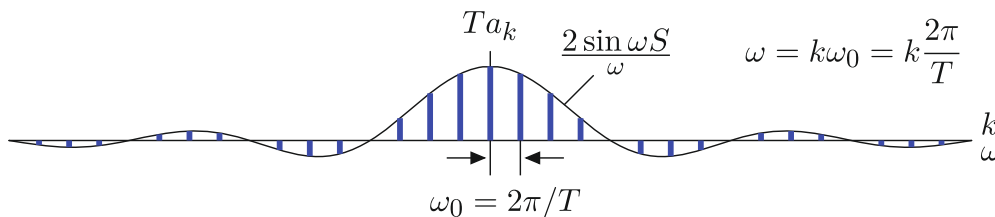
تبدیل فوریه

FOURIER TRANSFORM

Doubling period doubles # of harmonics in given frequency interval.



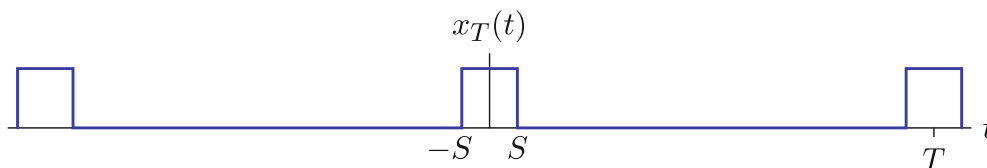
$$a_k = \frac{1}{T} \int_{-T/2}^{T/2} x_T(t) e^{-j \frac{2\pi}{T} kt} dt = \frac{1}{T} \int_{-S}^S e^{-j \frac{2\pi}{T} kt} dt = \frac{\sin \frac{2\pi k S}{T}}{\pi k} = \frac{2}{T} \frac{\sin \omega S}{\omega}$$



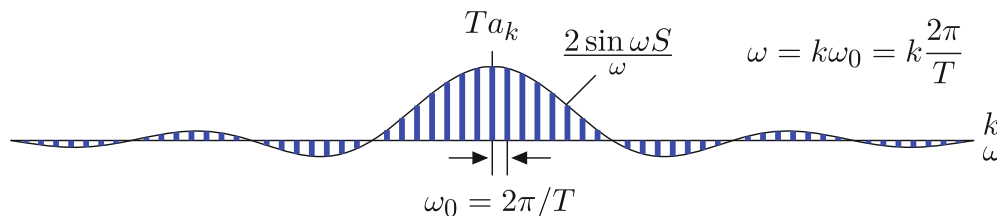
تبدیل فوریه

FOURIER TRANSFORM

As $T \rightarrow \infty$, discrete harmonic amplitudes $\rightarrow$ a continuum $E(\omega)$.



$$a_k = \frac{1}{T} \int_{-T/2}^{T/2} x_T(t) e^{-j \frac{2\pi}{T} kt} dt = \frac{1}{T} \int_{-S}^S e^{-j \frac{2\pi}{T} kt} dt = \frac{\sin \frac{2\pi k S}{T}}{\pi k} = \frac{2}{T} \frac{\sin \omega S}{\omega}$$

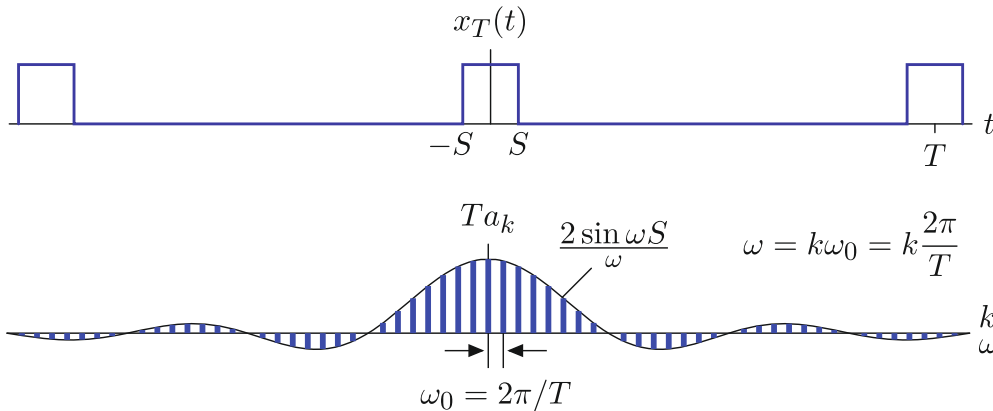


$$\lim_{T \rightarrow \infty} T a_k = \lim_{T \rightarrow \infty} \int_{-T/2}^{T/2} x(t) e^{-j \omega t} dt = \frac{2}{\omega} \sin \omega S = E(\omega)$$

تبدیل فوریه

FOURIER TRANSFORM

As $T \rightarrow \infty$, synthesis sum $\rightarrow$ integral.



$$\lim_{T \rightarrow \infty} T a_k = \lim_{T \rightarrow \infty} \int_{-T/2}^{T/2} x(t) e^{-j\omega t} dt = \frac{2}{\omega} \sin \omega S = E(\omega)$$

$$x(t) = \sum_{k=-\infty}^{\infty} \underbrace{\frac{1}{T} E(\omega)}_{a_k} e^{j\frac{2\pi}{T} kt} = \sum_{k=-\infty}^{\infty} \frac{\omega_0}{2\pi} E(\omega) e^{j\omega t} \rightarrow \frac{1}{2\pi} \int_{-\infty}^{\infty} E(\omega) e^{j\omega t} d\omega$$

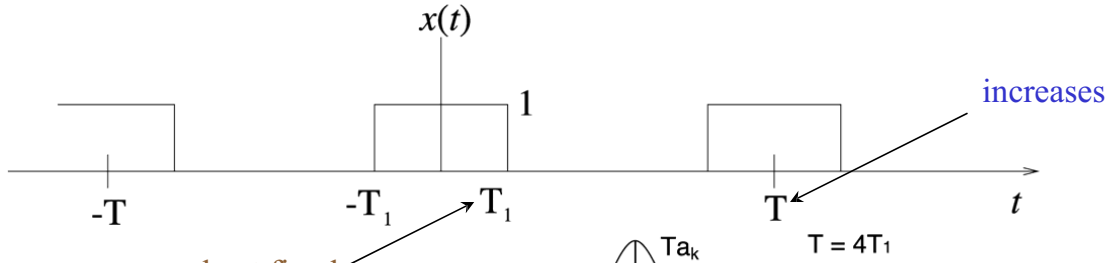
Fourier's Derivation of the CT Fourier Transform

- $x(t)$ - an aperiodic signal
 - view it as the limit of a periodic signal as $T \rightarrow \infty$
- For a periodic signal, the harmonic components are spaced $\omega_0 = 2\pi/T$ apart ...
- As $T \rightarrow \infty$, $\omega_0 \rightarrow 0$, and harmonic components are spaced closer and closer in frequency

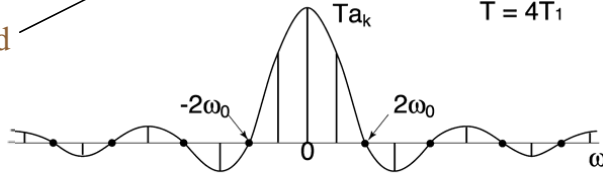


Fourier series $\longrightarrow$ Fourier integral

Motivating Example: Square wave

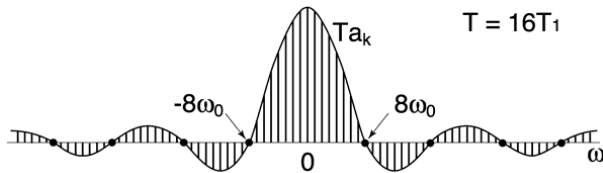
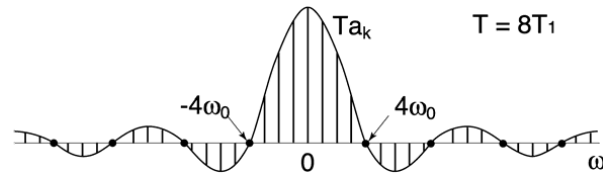


$$a_k = \frac{2 \sin(k\omega_0 T_1)}{k\omega_0 T}$$



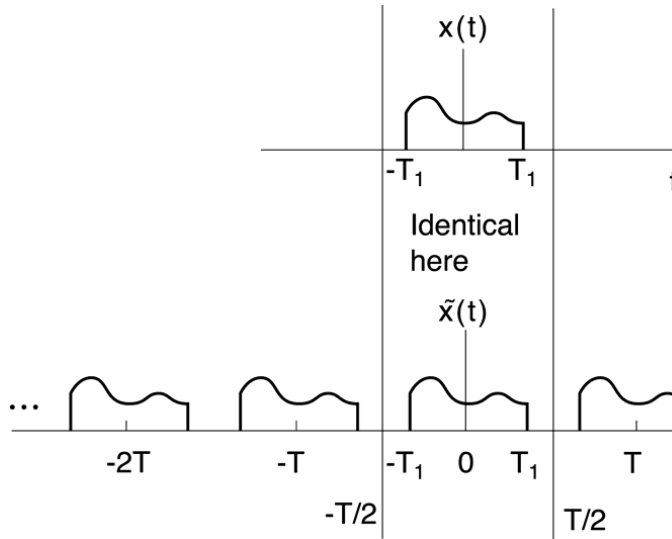
$\Downarrow$

$$Ta_k = \left. \frac{2 \sin \omega T_1}{\omega} \right|_{\omega = k\omega_0}$$



Discrete frequency points become denser in ω as T increases

So, on with the derivation ...



For simplicity, assume
 $x(t)$ a finite duration has (

$$\tilde{x}(t) = \begin{cases} x(t), & -\frac{T}{2} < t < \frac{T}{2} \\ \text{periodic}, & |t| > \frac{T}{2} \end{cases}$$

As $T \rightarrow \infty$, $\tilde{x}(t) = x(t)$ for all t

Derivation (continued)

$$\tilde{x}(t) = \sum_{k=-\infty}^{\infty} a_k e^{jk\omega_0 t} \quad \left(\omega_0 = \frac{2\pi}{T} \right)$$

$$a_k = \frac{1}{T} \int_{-\frac{T}{2}}^{\frac{T}{2}} \tilde{x}(t) e^{-jk\omega_0 t} dt = \frac{1}{T} \int_{-\frac{T}{2}}^{\frac{T}{2}} x(t) e^{-jk\omega_0 t} dt$$

$\uparrow$
 $\tilde{x}(t) = x(t)$ in this interval

$$= \frac{1}{T} \int_{-\infty}^{\infty} x(t) e^{-jk\omega_0 t} dt \quad (1)$$

If we define

$$X(j\omega) = \int_{-\infty}^{\infty} x(t) e^{-j\omega t} dt$$

then Eq.(1) $\Rightarrow$

$$a_k = \frac{X(jk\omega_0)}{T}$$

Derivation (continued)

Thus, for $-\frac{T}{2} < t < \frac{T}{2}$

$$\begin{aligned}x(t) = \tilde{x}(t) &= \sum_{k=-\infty}^{\infty} \underbrace{\frac{1}{T} X(jk\omega_0)}_{a_k} e^{jk\omega_0 t} \\&= \frac{1}{2\pi} \sum_{k=-\infty}^{\infty} \omega_0 X(jk\omega_0) e^{jk\omega_0 t} \\&\quad \Downarrow\end{aligned}$$

As $T \rightarrow \infty$, $\sum \omega_0 \rightarrow \int d\omega$, we get the CT Fourier Transform pair

$$x(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} X(j\omega) e^{j\omega t} d\omega \quad \text{Synthesis equation}$$

$$X(j\omega) = \int_{-\infty}^{\infty} x(t) e^{-j\omega t} dt \quad \text{Analysis equation}$$

For what kinds of signals can we do this?

(1) It works also even if $x(t)$ is infinite duration, but satisfies:

a) Finite energy $\int_{-\infty}^{\infty} |x(t)|^2 dt < \infty$

In this case, there is *zero* energy in the error

$$e(t) = x(t) - \frac{1}{2\pi} \int_{-\infty}^{\infty} X(j\omega) e^{j\omega t} d\omega \quad \text{Then} \quad \int_{-\infty}^{\infty} |e(t)|^2 dt = 0$$

b) Dirichlet conditions

(i) $\frac{1}{2\pi} \int_{-\infty}^{\infty} X(j\omega) e^{j\omega t} d\omega = x(t)$ at points of continuity

(ii) $\frac{1}{2\pi} \int_{-\infty}^{\infty} X(j\omega) e^{j\omega t} d\omega = \text{midpoint at discontinuity}$

(iii) Gibb's phenomenon

c) By allowing impulses in $x(t)$ or in $X(j\omega)$, we can represent even *more* signals

E.g. It allows us to consider *FT* for *periodic* signals

تبدیل فوری‌ی پیوسته-زمان (۱)

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مثال‌هایی از
تبدیل‌های
فوری‌ه

Example #1

(a) $x(t) = \delta(t)$

$$X(j\omega) = \int_{-\infty}^{\infty} \delta(t) e^{-j\omega t} dt = 1$$

$\Downarrow$

$$\delta(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{j\omega t} d\omega$$

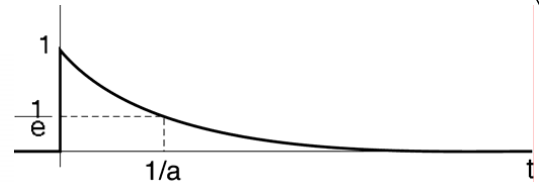
— Synthesis equation for $\delta(t)$

(b) $x(t) = \delta(t - t_0)$

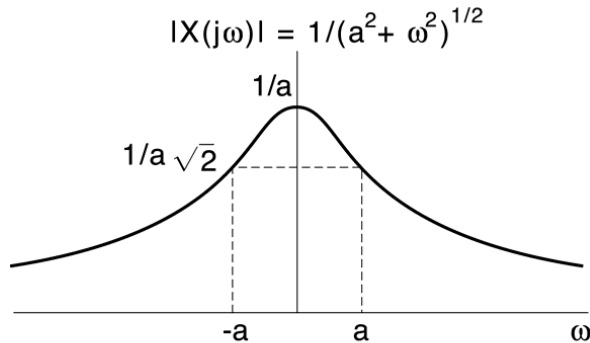
$$\begin{aligned} X(j\omega) &= \int_{-\infty}^{\infty} \delta(t - t_0) e^{-j\omega t} dt \\ &= e^{-j\omega t_0} \end{aligned}$$

Example #2: Exponential function

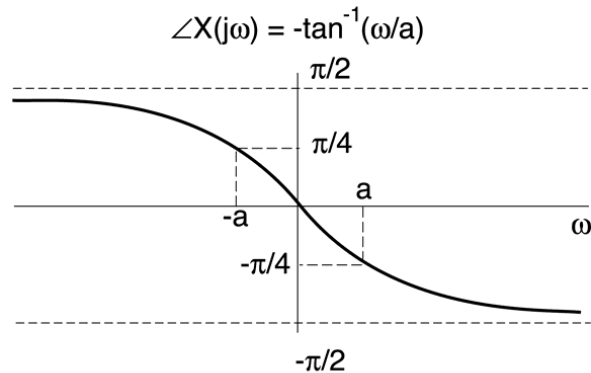
$$x(t) = e^{-at}u(t), a > 0$$



$$\begin{aligned} X(j\omega) &= \int_{-\infty}^{\infty} x(t)e^{-j\omega t} dt = \int_0^{\infty} \underbrace{e^{-at}e^{-j\omega t}}_{e^{-(a+j\omega)t}} dt \\ &= -\left(\frac{1}{a+j\omega}\right) e^{-(a+j\omega)t} \bigg|_0^{\infty} = \frac{1}{a+j\omega} \end{aligned}$$

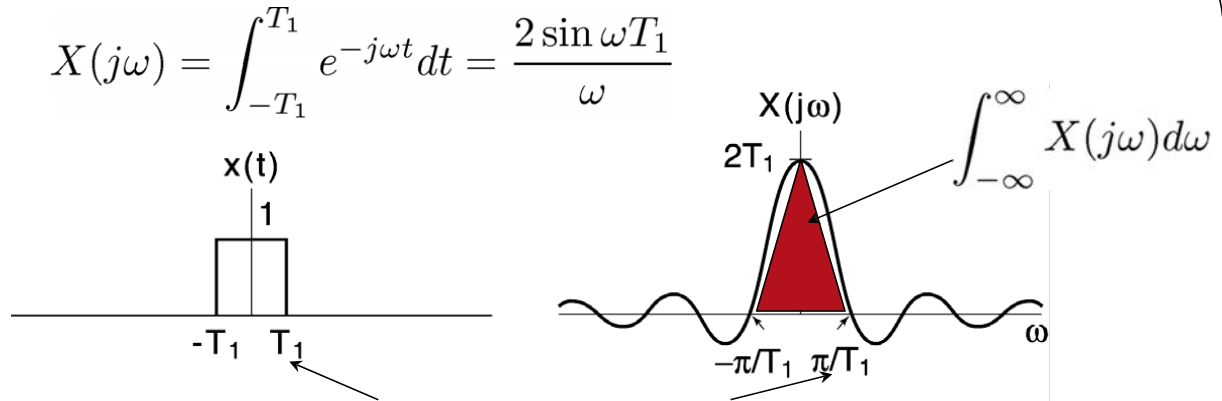


Even symmetry



Odd symmetry

Example #3: A square pulse in the time-domain



Note the inverse relation between the two widths $\Rightarrow$ **Uncertainty principle**

Useful facts about CTFT's

$$X(0) = \int_{-\infty}^{\infty} x(t) dt$$

Example above: $\int_{-\infty}^{\infty} x(t) dt = 2T_1 = X(0)$

$$x(0) = \frac{1}{2\pi} \int_{-\infty}^{\infty} X(j\omega) d\omega$$

Ex. above: $x(0) = 1 = \frac{1}{2\pi} \int_{-\infty}^{\infty} X(j\omega) d\omega$

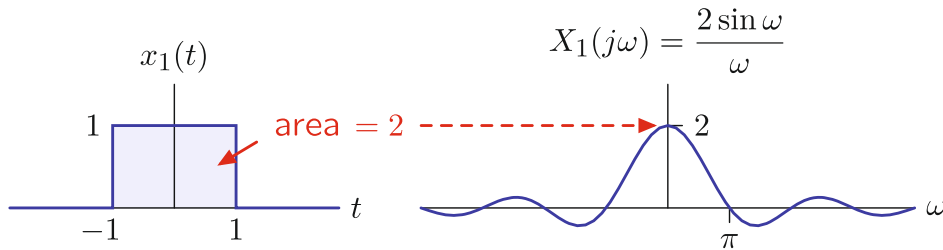
$$= \frac{1}{2\pi} \times (\text{Area of the triangle})$$

خصوصیات تبدیل فوری-زمان

رابطه‌ی حوزه‌ی زمان و حوزه‌ی فرکانس (۱)

The value of $X(j\omega)$ at $\omega = 0$ is the integral of $x(t)$ over time t .

$$X(j\omega)|_{\omega=0} = \int_{-\infty}^{\infty} x(t)e^{-j\omega t} dt = \int_{-\infty}^{\infty} x(t)e^{j0t} dt = \int_{-\infty}^{\infty} x(t) dt$$

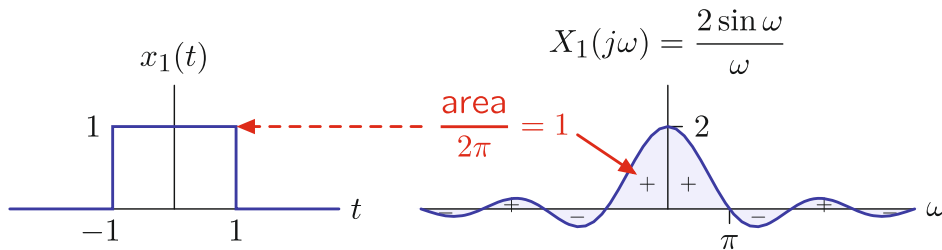


خصوصیات تبدیل فوریه‌ی پیوسته-زمان

رابطه‌ی حوزه‌ی زمان و حوزه‌ی فرکانس (۲)

The value of $x(0)$ is the integral of $X(j\omega)$ divided by 2π .

$$x(0) = \frac{1}{2\pi} \int_{-\infty}^{\infty} X(j\omega) e^{j\omega t} d\omega = \frac{1}{2\pi} \int_{-\infty}^{\infty} X(j\omega) d\omega$$

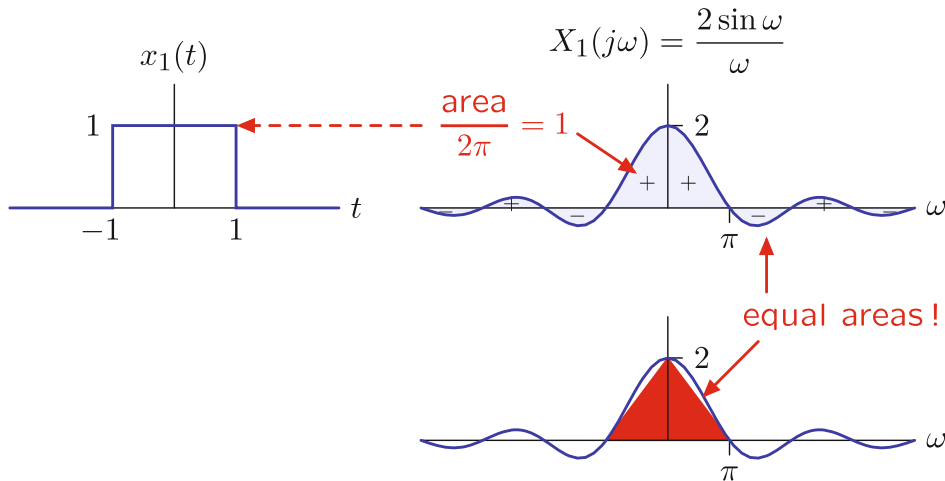


خصوصیات تبدیل فوریه پیوسته-زمان

رابطه‌ی حوزه‌ی زمان و حوزه‌ی فرکانس (۳)

The value of $x(0)$ is the integral of $X(j\omega)$ divided by 2π .

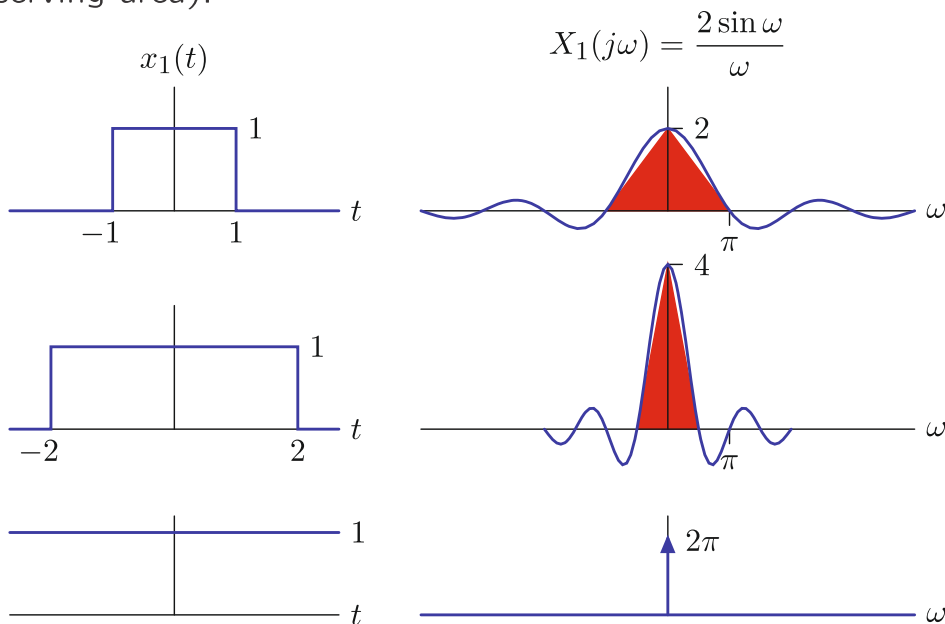
$$x(0) = \frac{1}{2\pi} \int_{-\infty}^{\infty} X(j\omega) e^{j\omega t} d\omega = \frac{1}{2\pi} \int_{-\infty}^{\infty} X(j\omega) d\omega$$



خصوصیات تبدیل فوریه‌ی پیوسته-زمان

رابطه‌ی حوزه‌ی زمان و حوزه‌ی فرکانس (۴)

Stretching time compresses frequency and increases amplitude (preserving area).

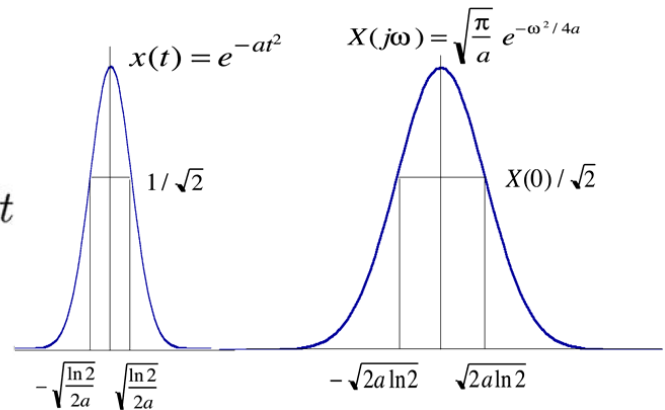


New way to think about an impulse!

Example #4: $x(t) = e^{-at^2}$ — A Gaussian, important in probability, optics, etc.

$$\begin{aligned}
 & X(j\omega) \\
 = & \int_{-\infty}^{\infty} e^{-at^2} e^{-j\omega t} dt \\
 = & \int_{-\infty}^{\infty} e^{-a \left[t^2 + j\frac{\omega}{a}t + \left(\frac{j\omega}{2a}\right)^2 \right] + a\left(\frac{j\omega}{2a}\right)^2} dt \\
 = & \underbrace{\left[\int_{-\infty}^{\infty} e^{-a\left(t + \frac{j\omega}{2a}\right)^2} dt \right]}_{\sqrt{\pi}/a} \cdot e^{-\frac{\omega^2}{4a}} \\
 = & \sqrt{\frac{\pi}{a}} e^{-\frac{\omega^2}{4a}}
 \end{aligned}$$

Also a Gaussian!



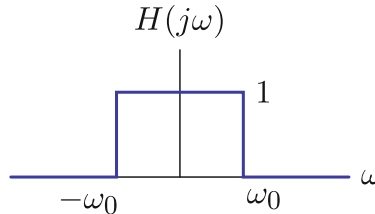
(Pulse width in t) • (Pulse width in ω)
 $\Rightarrow \Delta t \cdot \Delta \omega \sim (1/a^{1/2}) \cdot (a^{1/2}) = 1$

Uncertainty Principle! Cannot make both Δt and $\Delta \omega$ arbitrarily small.

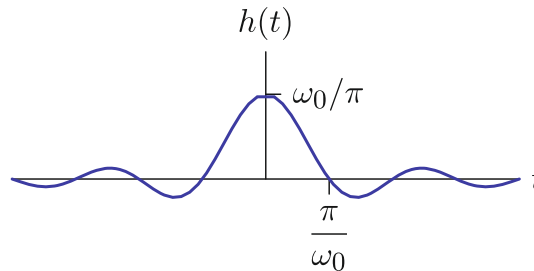
تبدیل وارون فوری

INVERSE FOURIER TRANSFORM

Find the impulse response of an “ideal” low pass filter.



$$h(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} H(j\omega) e^{j\omega t} d\omega = \frac{1}{2\pi} \int_{-\omega_0}^{\omega_0} e^{j\omega t} d\omega = \frac{1}{2\pi} \left. \frac{e^{j\omega t}}{jt} \right|_{-\omega_0}^{\omega_0} = \frac{\sin \omega_0 t}{\pi t}$$



This result is not so easily obtained without inverse relation.

خاصیت دوغانی

تبدیل مستقیم و تبدیل وارون فوری

DUALITY

The Fourier transform and its inverse have very similar forms.

$$X(j\omega) = \int_{-\infty}^{\infty} x(t) e^{-j\omega t} dt \quad (\text{Fourier transform})$$

$$x(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} X(j\omega) e^{j\omega t} d\omega \quad (\text{"inverse" Fourier transform})$$

Convert one to the other by

- $t \rightarrow \omega$
- $\omega \rightarrow -t$
- scale by 2π

خاصیت دوغانی

تبدیل مستقیم و تبدیل وارون فوری

DUALITY

The Fourier transform and its inverse have very similar forms.

$$X(j\omega) = \int_{-\infty}^{\infty} x(t)e^{-j\omega t} dt$$

$$x(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} X(j\omega)e^{j\omega t} d\omega$$

Two differences:

- minus sign: flips time axis (or equivalently, frequency axis)
- divide by 2π (or multiply in the other direction)

$$x_1(t) = f(t) \leftrightarrow X_1(j\omega) = g(\omega)$$



$$\omega \rightarrow t \quad t \rightarrow \omega ; \text{ flip } ; \times 2\pi$$

$$x_2(t) = g(t) \leftrightarrow X_2(j\omega) = 2\pi f(-\omega)$$

خاصیت دوگانی

تبدیل مستقیم و تبدیل وارون فوری

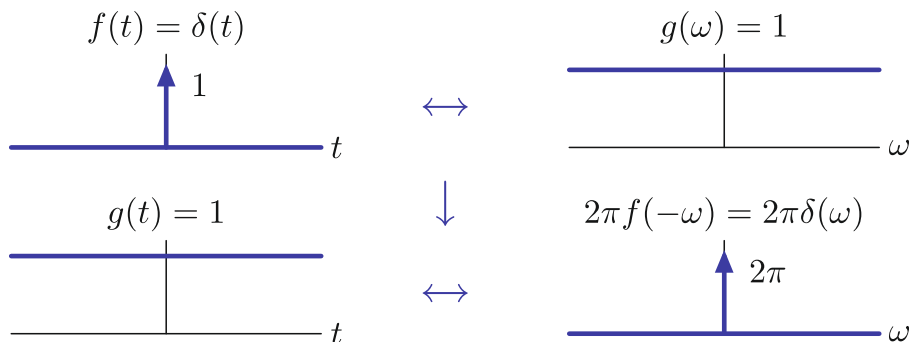
DUALITY

Using duality to find new transform pairs.

$$x_1(t) = f(t) \leftrightarrow X_1(j\omega) = g(\omega)$$

$$\omega \rightarrow t \quad t \rightarrow \omega ; \text{ flip ; } \times 2\pi$$

$$x_2(t) = g(t) \leftrightarrow X_2(j\omega) = 2\pi f(-\omega)$$



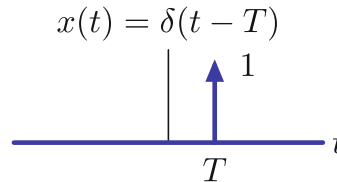
The function $g(t) = 1$ does not have a Laplace transform!

خاصیت دوگانی

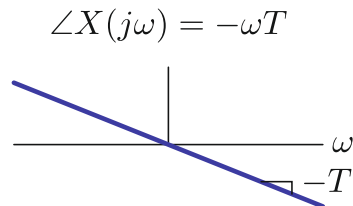
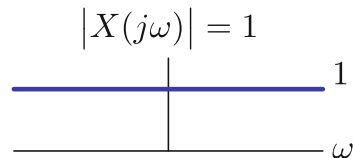
تبدیل مستقیم و تبدیل وارون فوری

DUALITY

Fourier transform of delayed impulse: $\delta(t - T) \leftrightarrow e^{-j\omega T}$.



$$X(j\omega) = \int_{-\infty}^{\infty} \delta(t - T) e^{-j\omega t} dt = e^{-j\omega T}$$



تبدیل فوریه‌ی پیوسته-زمان (۱)

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تبدیل‌های
فوریه‌ی
سیگنال‌های
متناوب

تبدیل فوری‌ی سیگنال‌های متناوب

تبدیل فوری‌ی سینوس

Using duality to find the Fourier transform of an eternal sinusoid.

$$\delta(t - T) \leftrightarrow e^{-j\omega T}$$

$$\omega \rightarrow t \quad t \rightarrow \omega ; \text{ flip } ; \times 2\pi$$

$$e^{-jtT} \leftrightarrow 2\pi\delta(\omega + T)$$

$T \rightarrow \omega_0 :$

$$e^{-j\omega_0 t} \leftrightarrow 2\pi\delta(\omega + \omega_0)$$

$$x(t) = x(t + T) = \sum_{k=-\infty}^{\infty} a_k e^{j\frac{2\pi}{T}kt} \quad \begin{matrix} \text{CTFS} \\ \longleftrightarrow \end{matrix} \quad \{a_k\}$$

$$x(t) = x(t + T) = \sum_{k=-\infty}^{\infty} a_k e^{j\frac{2\pi}{T}kt} \quad \begin{matrix} \text{CTFT} \\ \longleftrightarrow \end{matrix} \quad \sum_{k=-\infty}^{\infty} 2\pi a_k \delta\left(\omega - \frac{2\pi}{T}k\right)$$

CT Fourier Transforms of Periodic Signals

Suppose

$$X(j\omega) = \delta(\omega - \omega_0)$$

$\Downarrow$

$$x(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \delta(\omega - \omega_0) e^{j\omega t} d\omega = \frac{1}{2\pi} e^{j\omega_0 t} \text{ — periodic in } t \text{ with frequency } \omega_0$$

That is

$$e^{j\omega_0 t} \leftrightarrow 2\pi\delta(\omega - \omega_0)$$

— All the energy is concentrated in one frequency — ω_0

More generally

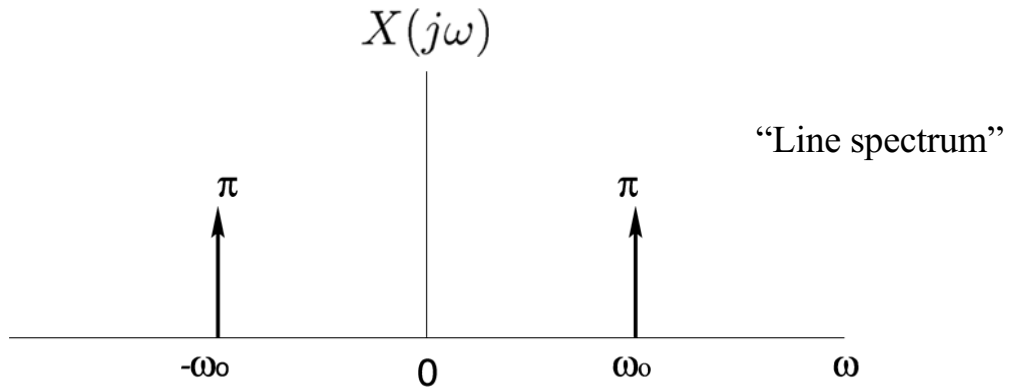
$$x(t) = \sum_{k=-\infty}^{\infty} a_k e^{jk\omega_0 t} \leftrightarrow X(j\omega) = \sum_{k=-\infty}^{\infty} 2\pi a_k \delta(\omega - k\omega_0)$$

Example #4:

$$x(t) = \cos \omega_0 t = \frac{1}{2}e^{j\omega_0 t} + \frac{1}{2}e^{-j\omega_0 t}$$

$\Downarrow$

$$X(j\omega) = \pi\delta(\omega - \omega_0) + \pi\delta(\omega + \omega_0)$$

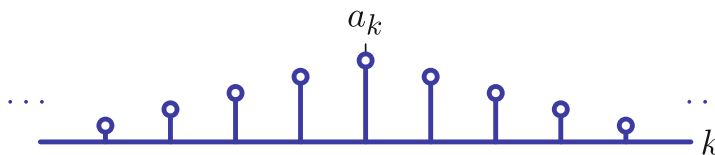
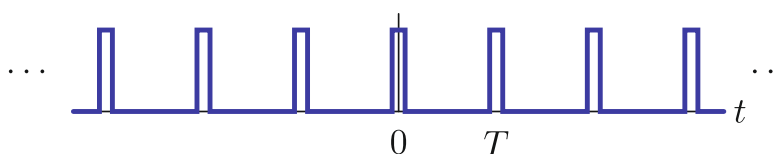


تبدیل فوری‌ی سیگنال‌های متناوب

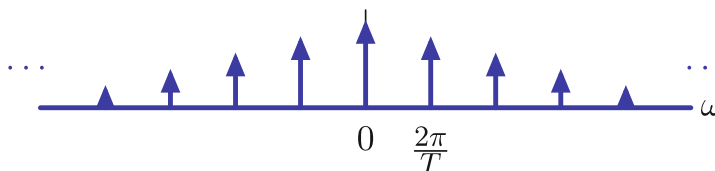
رابطه‌ی کلی

Each term in the Fourier series is replaced by an impulse.

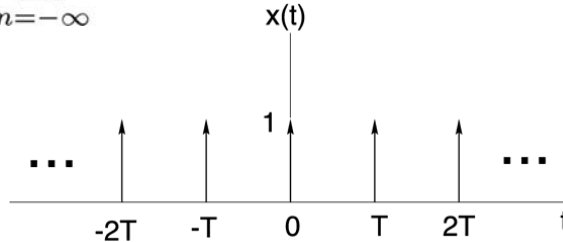
$$x(t) = \sum_{k=-\infty}^{\infty} x_p(t - kT)$$



$$X(j\omega) = \sum_{k=-\infty}^{\infty} 2\pi a_k \delta\left(\omega - k\frac{2\pi}{T}\right)$$



Example #5: $x(t) = \sum_{n=-\infty}^{\infty} \delta(t - nT)$ — Sampling function

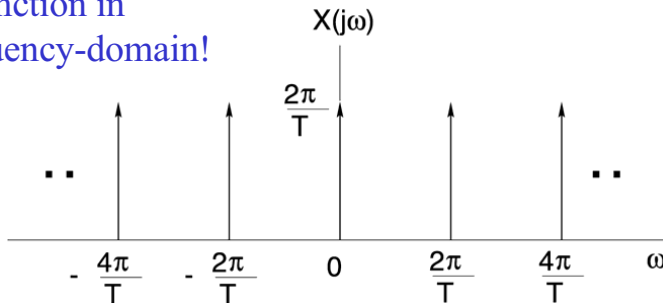


$$x(t) \leftrightarrow a_k = \frac{1}{T} \int_{-T/2}^{T/2} x(t) e^{-jk\omega_0 t} dt = \frac{1}{T}$$

$\Downarrow$

$$X(j\omega) = \sum_{n=-\infty}^{\infty} \underbrace{\frac{2\pi}{T}}_{2\pi a_k} \delta\left(\omega - \underbrace{\frac{k2\pi}{T}}_{k\omega_0}\right)$$

Same function in
the frequency-domain!



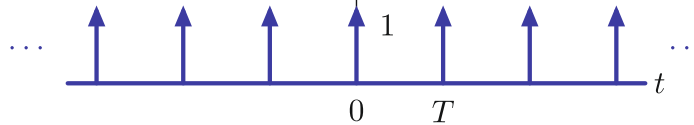
Note: (period in t) T
 $\Leftrightarrow$ (period in ω) $2\pi/T$
 Inverse relationship again!

تبدیل فوری‌ی سیگنال‌های متناوب

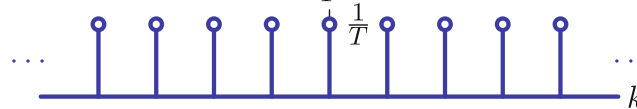
تبدیل فوری‌ی قطار ضربه

The Fourier transform of an impulse train is an impulse train.

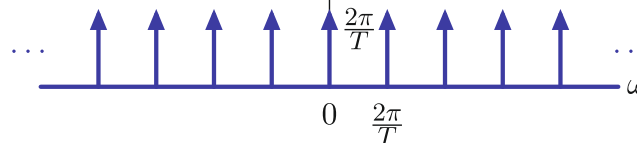
$$x(t) = \sum_{k=-\infty}^{\infty} \delta(t - kT)$$



$$a_k = \frac{1}{T} \quad \forall k$$



$$X(j\omega) = \sum_{k=-\infty}^{\infty} \frac{2\pi}{T} \delta(\omega - k\frac{2\pi}{T})$$



تبدیل فوریه‌ی پیوسته-زمان (۱)

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خصوصیات
تبدیل فوریه‌ی
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Properties of the CT Fourier Transform

1) Linearity $ax(t) + by(t) \leftrightarrow aX(j\omega) + bY(j\omega)$

2) Time Shifting $x(t - t_0) \leftrightarrow e^{-j\omega t_0} X(j\omega)$

Proof:
$$\int_{-\infty}^{\infty} \underbrace{x(t - t_0)}_{t'} e^{-j\omega t} dt = e^{-j\omega t_0} \underbrace{\int_{-\infty}^{\infty} x(t') e^{-j\omega t'} dt'}_{X(j\omega)}$$

FT magnitude unchanged

$$|e^{-j\omega t_0} X(j\omega)| = |X(j\omega)|$$

Linear change in *FT* phase

$$\angle(e^{-j\omega t_0} X(j\omega)) = \angle X(j\omega) - \omega t_0$$

Properties (continued)

3) Conjugate Symmetry

$$x(t) \text{ real} \leftrightarrow X(-j\omega) = X^*(j\omega)$$

$\Downarrow$

$$|X(-j\omega)| = |X(j\omega)| \quad \text{Even}$$

$$\angle X(-j\omega) = -\angle X(j\omega) \quad \text{Odd}$$

$$\text{Re}\{X(-j\omega)\} = \text{Re}\{X(j\omega)\} \quad \text{Even}$$

$$\text{Im}\{X(-j\omega)\} = -\text{Im}\{X(j\omega)\} \quad \text{Odd}$$

The Properties Keep on Coming ...

4) Time-Scaling $x(at) \longleftrightarrow \frac{1}{|a|} X\left(j\frac{\omega}{a}\right)$ E.g. $a > 1 \rightarrow at > t$
 $\Downarrow a = -1$ compressed in time $\Leftrightarrow$
 stretched in frequency
 $x(-t) \longleftrightarrow X(-j\omega)$

a) $x(t)$ real and even $\Downarrow$ $x(t) = x(-t)$
 $\Rightarrow X(j\omega) = X(-j\omega) = X^*(j\omega)$ - Real & even

b) $x(t)$ real and odd $x(t) = -x(-t)$
 $\Rightarrow X(j\omega) = -X(-j\omega) = -X^*(j\omega)$ - Purely imaginary & odd

c)

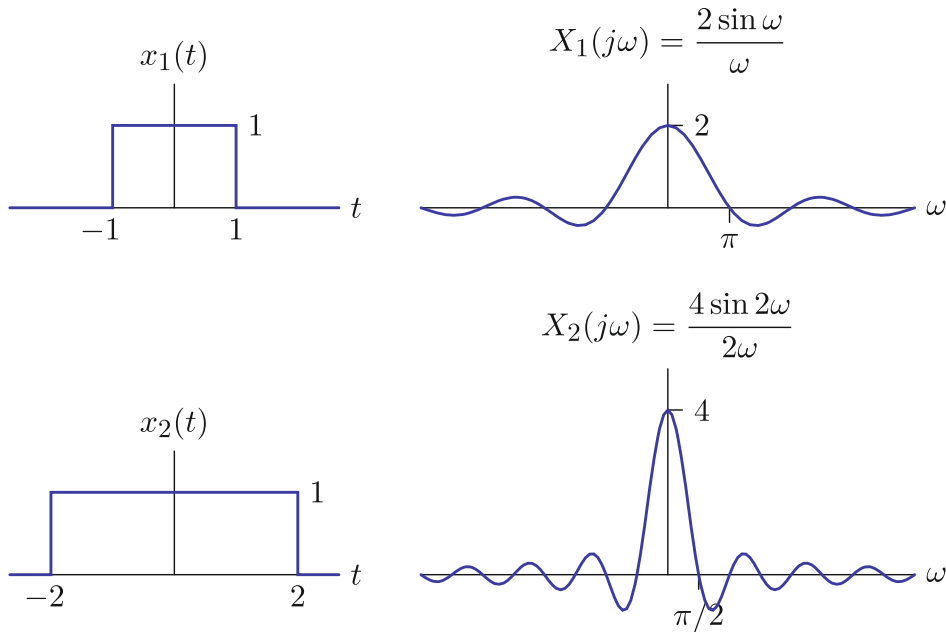
$$X(j\omega) = \text{Re}\{X(j\omega)\} + j\text{Im}\{X(j\omega)\}$$

$$\text{For real } x(t) = \underset{\uparrow}{\text{Ev}\{x(t)\}} + \underset{\uparrow}{j\text{Od}\{x(t)\}}$$

خصوصیات تبدیل فوریه‌ی پیوسته-زمان

کشش زمانی موجب فشردگی فرکانسی می‌شود

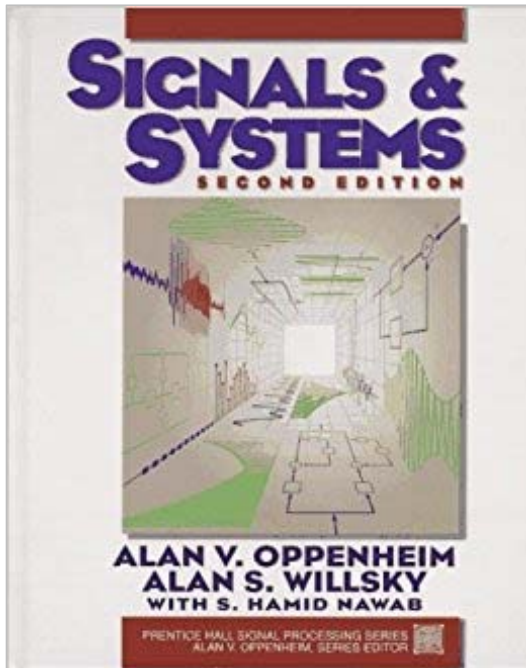
Stretching time compresses frequency.



تبدیل فوریه‌ی پیوسته-زمان (۱)

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منابع



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Chapter 4